

# PHYS 320/420: Lecture 20

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$$E(t) = \langle E \rangle_t = \sum_n P_n(t) E_n$$

$$\frac{dE}{dt} = \sum_{nm} J_{nm} E_n \quad (*)$$

$$= \frac{1}{2} \sum_{nm} J_{nm} (E_n - E_m)$$

example:

$$J_{nm} = \Omega_{nm} \rho_m - \Omega_{mn} \rho_n$$



$$J_{nn} = 0, \quad J_{nm} = -J_{mn}$$

$$(*) = J_{12} E_1 + J_{21} E_2 + J_{23} E_2 + J_{32} E_3$$

$$(**) = J_{12} (E_1 - E_2) + J_{23} (E_2 - E_3)$$

using  $J_{21} = -J_{12}$ ,  $J_{32} = -J_{23}$

$$(***) = J_{21} (E_2 - E_1) + J_{32} (E_3 - E_2)$$

$$(*) = \frac{1}{2} [(**) + (***)] = \frac{1}{2} \sum_{nm} J_{nm} (E_n - E_m)$$

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# General definitions:

state  $(n)$  has property  $A_n$   
(energy, etc.)

avg. of that property in sys at time  $t$   $= A(t) \equiv \sum_n P_n(t) A_n$

$$\dot{A}(t) = \frac{dA}{dt} = \frac{1}{2} \sum_{nm} J_{nm}(t) (A_n - A_m)$$

transition  $(m) \rightarrow (n)$  has property  $B_{nm}$

"production rate"  $\dot{B}(t) \equiv \frac{1}{2} \sum_{nm} J_{nm}(t) B_{nm}$

abuse of notation: sometimes

$\dot{B}(t)$  exists,  $\nexists$  no  $B(t)$

exception: special case

where  $B_{nm} = A_n - A_m$

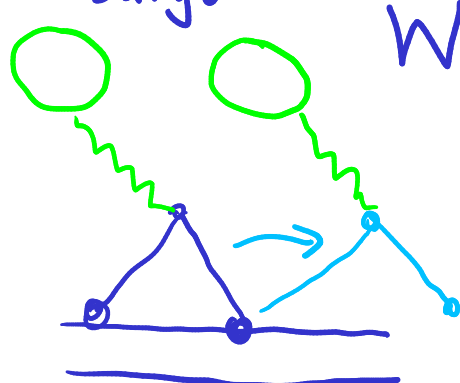
trans. property      diff. in state properties

$$\Rightarrow \dot{B}(t) = \frac{dA}{dt}$$

where  
 $A(t) = \sum_n P_n(t) A_n$   
 is well defined

example: work done by sys  
 during  $m \rightarrow n$  transition

motor  
 protein doing  
 work on  
 cargo

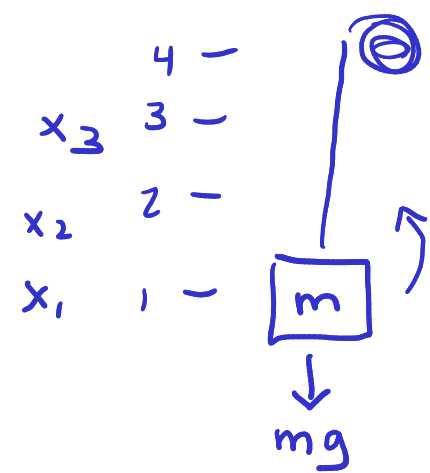


$$W_{nm}$$

$$\dot{W}(t) = \frac{1}{2} \sum_{nm} J_{nm} W_{nm}$$

= rate<sup>mean</sup> of work  
 done by sys

= power output of sys



$$\begin{aligned} W_{21} &= mg(x_2 - x_1) \\ &= mgx_2 - mgx_1 \\ &= U_2 - U_1 \end{aligned}$$

Let's define another transition property associated w/ any allowable trans. ( $\Omega_{nm} \neq 0 \Leftrightarrow \Omega_{mn} \neq 0$ ):

define "irreversibility"  $I_{nm}(t)$

$$\equiv k_B \ln \frac{\Omega_{nm} P_m(t)}{\Omega_{mn} P_n(t)}$$

by  
convention

$$= k_B \ln \frac{\text{avg. \# of } m \rightarrow n \text{ jumps / time}}{\text{avg. \# of } n \rightarrow m \text{ jumps / time}}$$

If "reverse" jumps  $n \rightarrow m$  are much less likely than "forward"  $m \rightarrow n$  jumps  
 $\Rightarrow I_{nm}(t)$  is large

production rate  $\dot{I}(t) = \frac{1}{2} \sum_{nm} J_{nm}(t) I_{nm}(t)$

$$= \frac{k_B}{2} \sum_{nm} (\Omega_{nm} P_m - \Omega_{mn} P_n) \ln \frac{\Omega_{nm} P_m}{\Omega_{mn} P_n}$$

every term  
has form:

$$(x - y) \ln \frac{x}{y}$$

$$= \begin{cases} > 0 & \text{if } x \neq y \\ = 0 & \text{if } x = y \end{cases}$$

$$\dot{I}(t) \geq 0$$

for us this is  
equiv. to 2nd law of  
thermodynamics (we will  
explain soon!)

valid at  
all times for  
all systems

(note: indep. of  
det. balance relation)

does not require stationary state

Note:  $\dot{I}(t) = 0$  can only happen

$$\text{if } \Omega_{nm} p_m(t) = \Omega_{mn} p_n(t)$$

for every connected  $(n, m)$

$$\Rightarrow J_{nm}(t) = 0 \text{ for every } (n, m)$$

This can only happen if the  
system is in equilibrium

$$\frac{dp_n}{dt} = \sum_m J_{nm}(t) \quad (\text{ESS})$$

What happens if detailed balance condition is also true?

$$\frac{\Omega_{nm}}{\Omega_{mn}} = e^{-\beta(E_n - E_m)} \quad \beta = \frac{1}{k_B T}$$

$$\dot{I} = \frac{k_B}{2} \sum_{nm} J_{nm} \ln \frac{\Omega_{nm}}{\Omega_{mn}} \frac{P_m}{P_n}$$

$$= \frac{-1}{2T} \sum_{nm} J_{nm} (E_n - E_m)$$

$$+ \frac{k_B}{2} \sum_{nm} J_{nm} \ln \frac{P_m}{P_n}$$

$$= \frac{-\dot{E}}{T} + \frac{1}{2} \sum_{nm} J_{nm} \left[ \underbrace{-k_B \ln p_n}_{S_n(t)} - \underbrace{(-k_B \ln p_m)}_{S_m(t)} \right]$$

= "entropy" of state n

$$\dot{I} = -\frac{\dot{E}}{T} + \dot{S} \quad \text{where} \quad \dot{S}(t) = \frac{d}{dt} S(t)$$

$$S(t) = \sum_n p_n(t) S_n(t)$$

$$= -k_B \sum_n p_n(t) \ln p_n(t)$$

"mean" entropy of system  
 $\Rightarrow$  typically called  
 entropy of system

define  $F(t) = E(t) - TS(t)$   
 = Helmholtz free energy

$$-T\dot{I} = \dot{E} - T\dot{S} = \dot{F}$$

$F(t)$  must have slope  $\leq 0$

b/c  $\dot{F} = -T\dot{I}$  and  $T > 0$   
 $\dot{I} \geq 0$

this free energy  
 must dec. or stay  
 constant over time